

# Optimal Static Potentials for Robust Macroscopic Quantum State Generation of Levitated Nanoparticles



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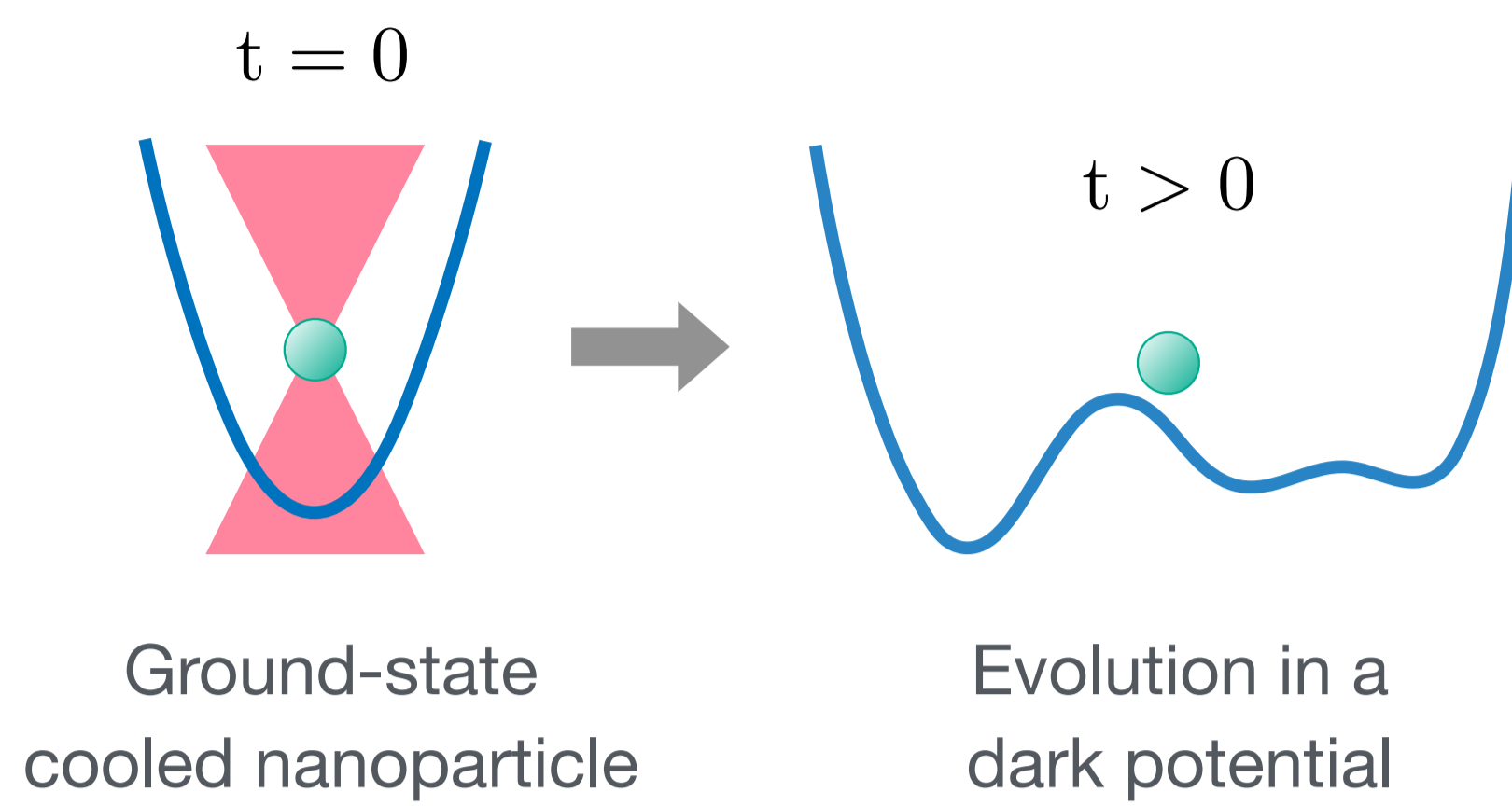
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## 1. Introduction

### Levitated nanoparticles in the quantum regime

- Ground state cooling achieved [1,2,3]
- Towards preparation of **macroscopic quantum superposition states**:
  - Coupling to nonlinear systems is challenging
  - Evolution in dark nonharmonic potentials [4,5]
  - Need for delocalization to enhance quantum features, which introduces decoherence



What is the **optimal potential shape**?

### Proposal

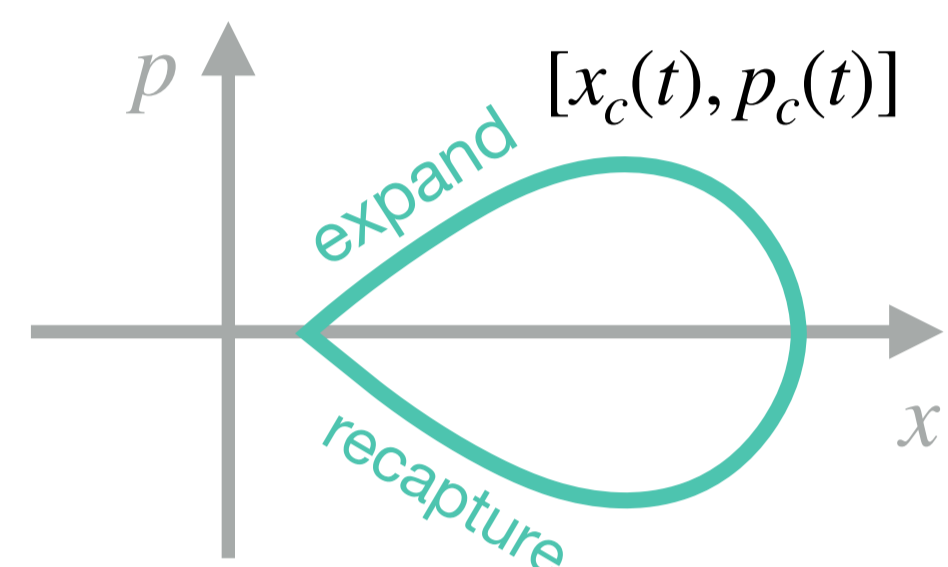
- Optimization of wide static potentials for the generation of macroscopic quantum states
- Optimization accounting for position-dependent **noise** sources within experimental setups
- Introduction of key **figures of merit** that allow for fast computation and capture macroscopicity and quantum signatures

## 2. Optimization

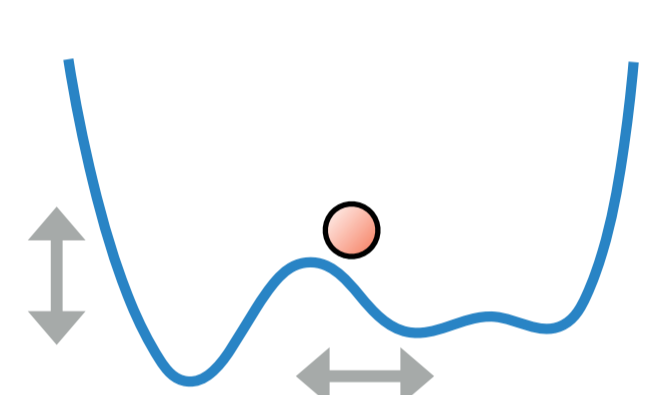
### Optimization of a static potential

$$V(x) = \frac{1}{2}m\omega^2 d^2 \sum_{n=1}^N \frac{c_n}{n!} \left(\frac{x}{d}\right)^n$$

- Constraints:
- Closed classical trajectory
  - Feasible potential
  - Fast run



### Optimization in the presence of position-dependent noise



- Fluctuations in position and amplitude of the potential

$$\tilde{V}(x, t) = V(x + x_0 \xi_1(t)) [1 + \xi_2(t)]$$

$$\langle \xi_i(t) \xi_j(t') \rangle = 2\pi S_i \delta_{ij} \delta(t - t')$$

### Figures of merit based on Gaussian dynamics

- Classical trajectory frame: phase-space shift onto a trajectory associated with the classical dynamics [6]

$$\hat{H} = \frac{\hat{p}^2}{2m} + V_s(\hat{x}) \rightarrow \hat{H}_c(t) = \frac{\hat{p}^2}{2m} + \frac{1}{2}m\omega^2 \sum_{n=2}^N \alpha_n(t) \hat{x}^n$$

$$\frac{\partial \hat{\rho}_c}{\partial t} = \frac{1}{i\hbar} [\hat{H}_c(t), \hat{\rho}_c(t)] - \frac{\Gamma_f(t)}{2X_\Omega^2} [\hat{x}, [\hat{x}, \hat{\rho}_c(t)]]$$

- Gaussian dynamics

$$\hat{H}_G(t) = \frac{\hat{p}^2}{2m} + \frac{1}{2}m\omega^2 \alpha_2(t) \hat{x}^2 \rightarrow \text{squeezing with parameter } \eta(t) \text{ and angle } \phi(t)$$

- First quantum correction

$$\hat{V}_{NG}(t) \simeq \frac{1}{2}m\omega^2 \alpha_3(t) \hat{x}^3 \rightarrow \hat{V}_{NG}^G(t) \simeq \beta_3(t) [\hat{x} \cos(\phi(t)) + \hat{p} \sin(\phi(t))]^3$$

- Figures of merit that allow for fast computation

Coherence length:  $\xi = \sqrt{8\langle \hat{x}_\phi^2 \rangle \mathcal{P}(t)}$   $\langle -\frac{\hat{x}_\phi}{2} | \hat{\rho}_G | -\frac{\hat{x}_\phi}{2} \rangle \propto \exp(-x^2 / \xi_\phi^2)$

- Cubicity and purity:

$$K(t) = \kappa_3(t) \mathcal{P}(t) \quad \kappa_3^2(t) = \left[ \int_0^t dt' \beta_3(t') \sin(\phi(t')) \right]^2 + \left[ \int_0^t dt' \beta_3(t') \cos(\phi(t')) \right]^2$$

## 4. Conclusion and Outlook

- We propose a new method for optimizing static potentials for the generation of macroscopic quantum states of levitated nanoparticles in the presence of noise.
- We introduce figures of merit that allow for fast computation and capture macroscopicity and quantum signatures.
- We obtain the optimal quartic potential for maximizing squeezing and for generating quantum cubic states.
- Outlook:**
  - Optimization within more general families of potentials
  - Optimal potentials for noise metrology

## 3. Example: quartic potentials

### Static quartic potential

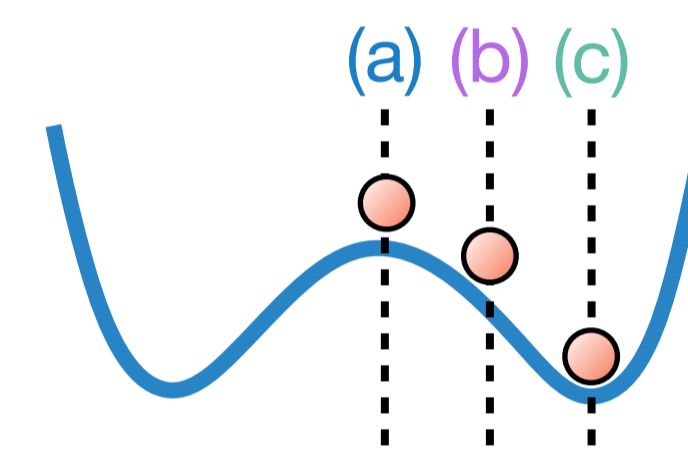
$$V(x) = \frac{1}{2}m\omega^2 \left( a(x - d_0)^2 + \frac{b}{2d^2}(x - d_0)^4 \right)$$

$a, b, d_0$   
optimization parameters

### Optimization of coherence length

- Optimal configurations:

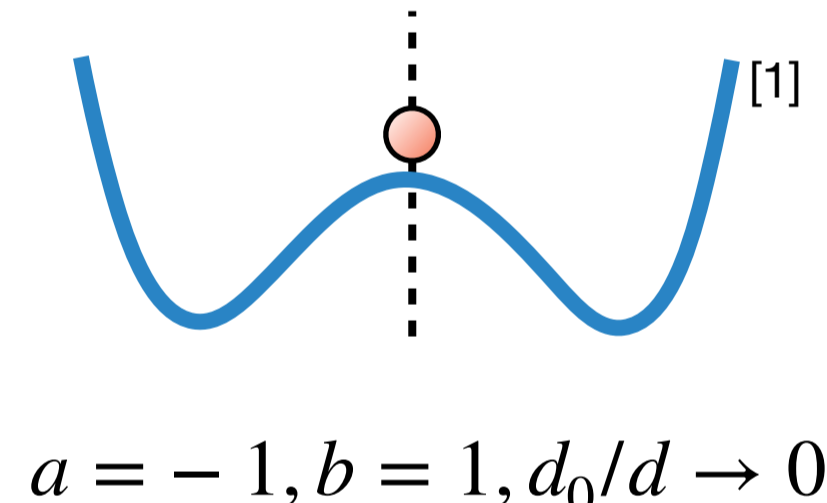
$$\Gamma_f(t) = \frac{2\pi X_\Omega^2}{\hbar^2} (S_1 X_\Omega^2 [V''(x_c(t))]^2 + S_2 [V'(x_c(t))]^2)$$



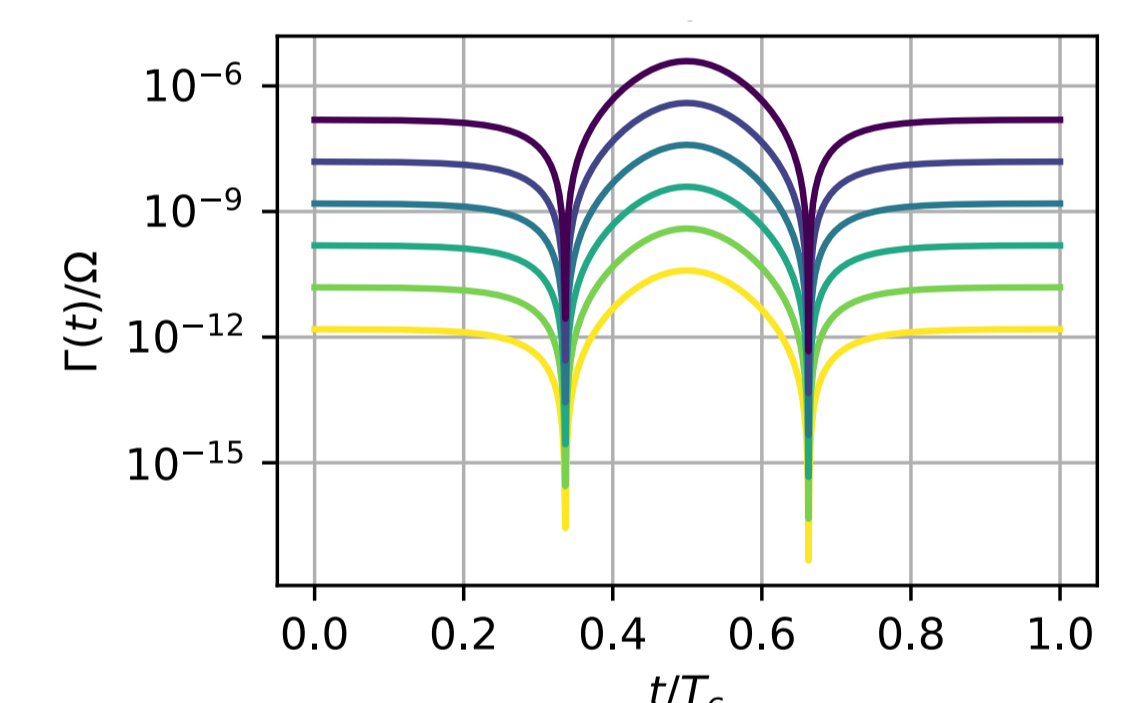
- (a) Small noise ( $S_i \omega \lesssim 1, i = 1, 2$ )
- (b) Large displacement noise ( $S_1 \omega \gtrsim 10$ )
- (c) Large intensity noise ( $S_2 \omega \gtrsim 10$ )

### Optimization of cubicity and purity

- Optimal potential:

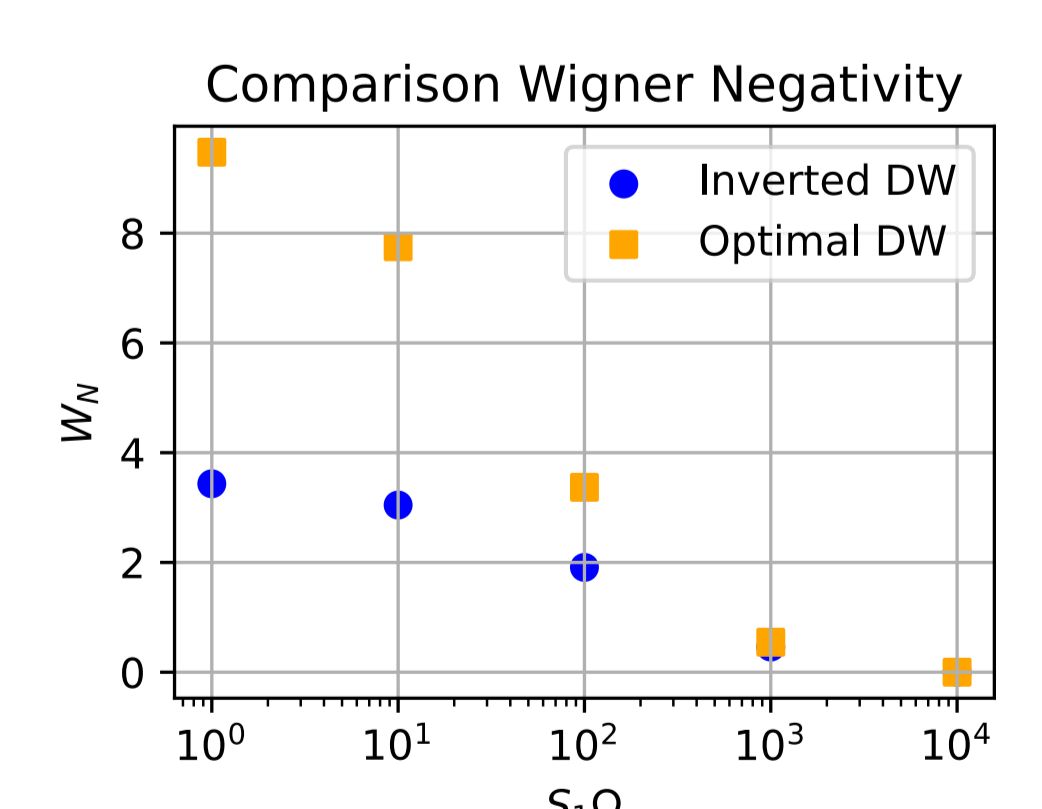
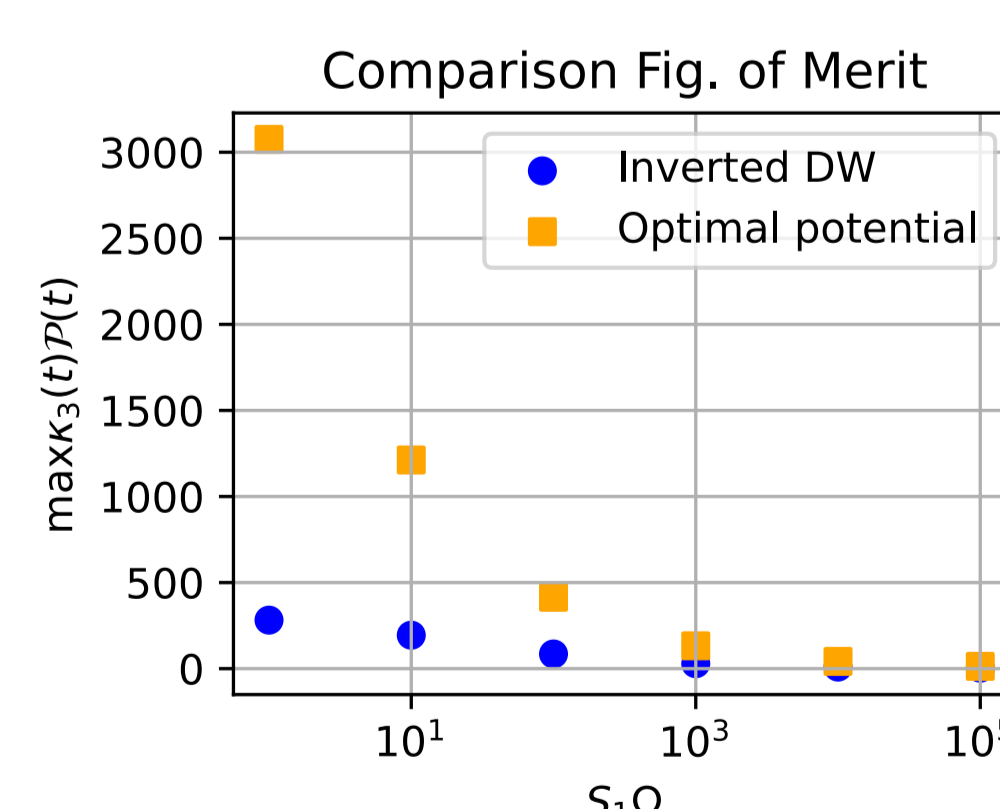
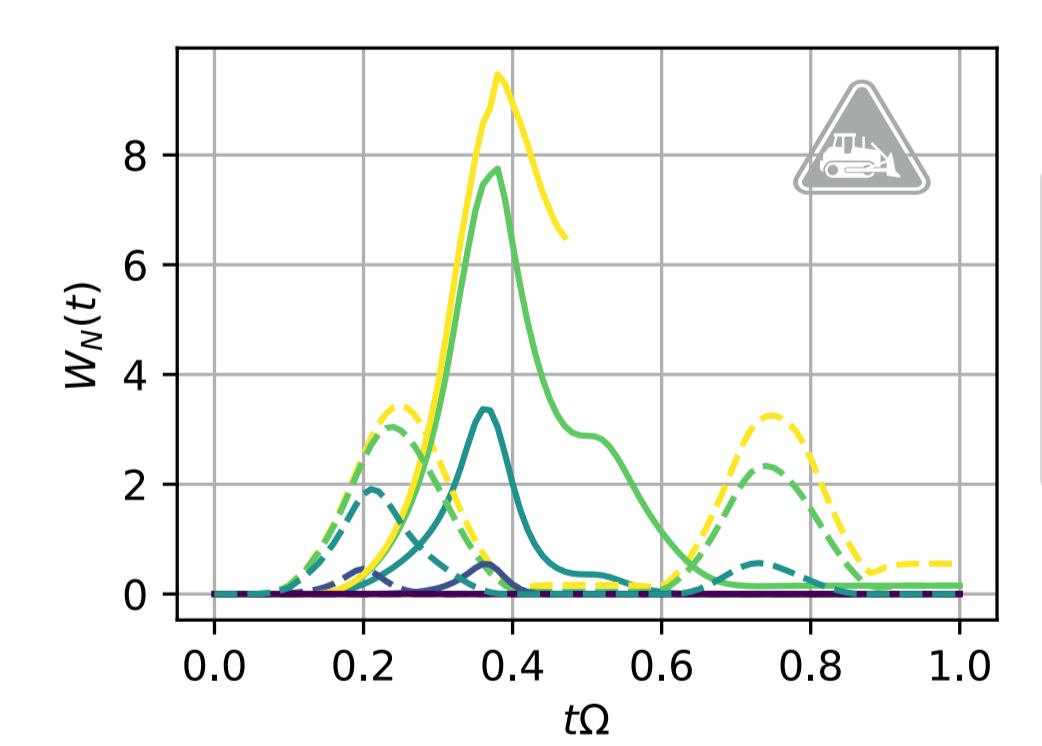
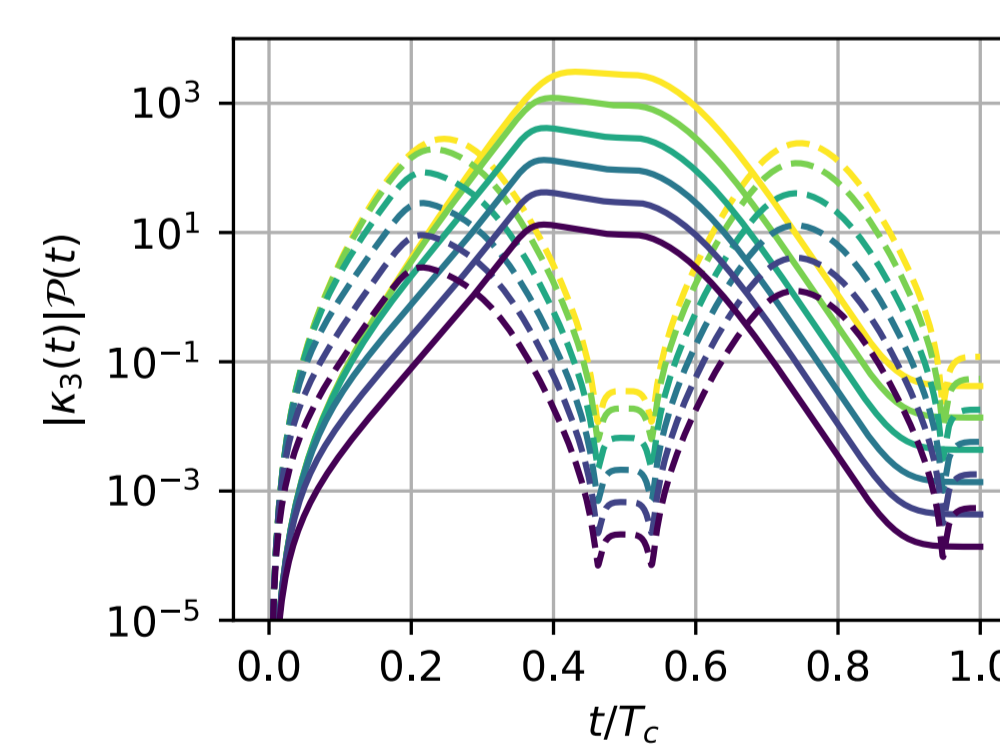


- Decoherence rate:



### Comparison with full quantum dynamics [7]

- Two potentials: DW and inverted DW
- Our figure of merit qualitatively captures the Wigner negativity features



## 5. References

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